Statistical Mechanics and Material Properties

By

Kunio TAKAHASHI
Tokyo Institute of Technology,
Tokyo 152-8552, JAPAN
Phone/Fax +81-3-5734-3915
takahak@ide.titech.ac.jp
http://www.ide.titech.ac.jp/~KT-lab/

© Only for the students in TAKAHASHI's class. All right reserved. Do not distribute to others. 2013

Statistical Mechanics and Thermodynamics

Thermodynamics

(0^{th} low: A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Interface essential to utilize thermodynamics



Grave-stone of Boltzmann...

Statistical Mechanics

Ergodic theorem : time average = space average, (in equilibrium)

You may have learnt Thermodynamics

Thermodynamics

(0th low: A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Everything is derived from these principles...

But!

- Definitions
 - Temperature
 - Free energy
 - Entropy
 - Equilibrium

Please remember when you have heard them first.

How they were defined?

Could you understand them as you could do in kinetics?

Definitions

Temperature

°C: water

°F: salt water

Empirical temp.

Temp. in science (physics or kinetics) is same as above?

Free energy

Helmholtz

$$F = U - TS$$

Gibbs

$$G = U + PV - TS$$

Free energy is same as the energy in kinetics?

Why they need to be defined as above?

Entropy

What is this?

Equilibrium

What is the definition?

You may have learnt Thermodynamics

Thermodynamics

(0th low: A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Everything is derived from these principles...

But!

- Definitions
 - Temperature
 - Free energy
 - Entropy
 - Equilibrium

Understanding of them (= statistical mechanics)

lead you to utilize

thermodynamics

Material properties

Let's define them step by step!

Temperature (first trial)

Example: Ideal gas (1)

Assumption

Number of molecule: N in volume : V

Velocity of molecule : c = (u, v, w)

probability: F(c) = f(u)f(v)f(w)

Momentum change = Pressure on area A

$$\int_{0}^{+\infty} \left(2mu \times f(u) \frac{N}{V} Audt\right) du = PAdt$$

$$PV = \int_{0}^{+\infty} (2mu \times f(u)Nu) du = mN \int_{-\infty}^{+\infty} u^{2} f(u) du = Nm \langle u^{2} \rangle$$

Comparison with eq. of state:

$$\left| \left\langle \frac{1}{2} mc^2 \right\rangle = \frac{3}{2} kT$$

Law of equi-partition of energy

$$c^{2} = u^{2} + v^{2} + w^{2}$$

$$\langle u^{2} \rangle = \langle v^{2} \rangle = \langle w^{2} \rangle$$

$$\langle u^{2} \rangle = \frac{1}{3} \langle c^{2} \rangle$$

$$\left\langle \frac{1}{2}mc^{2}\right\rangle = \frac{3}{2}kT$$

$$PV = Nm\frac{1}{3}\left\langle c^{2}\right\rangle = N\frac{2}{3}\left\langle \frac{1}{2}mc^{2}\right\rangle$$

PV = NkT

$$\left\langle \frac{1}{2}mu^2 \right\rangle = \left\langle \frac{1}{2}mv^2 \right\rangle = \left\langle \frac{1}{2}mw^2 \right\rangle = \frac{1}{2}kT$$

Example: Ideal gas (2)

$$c^2 = u^2 + v^2 + w^2$$

Probability function

$$F(c) = f(u)f(v)f(w)$$
 taking log: $\ln F(c) = \ln f(u) + \ln f(v) + \ln f(w)$

diff. with
$$u$$
:
$$\frac{d \ln F}{du} = \frac{d \ln f}{du} \Rightarrow \frac{d \ln F}{dc} \frac{dc}{du} = \frac{d \ln F}{dc} \frac{u}{c} = \frac{d \ln f}{du}$$

Const. independent of
$$c, u, v, w$$
:
$$\frac{1}{c} \frac{d \ln F}{dc} = \frac{1}{u} \frac{d \ln f}{du} = \frac{1}{v} \frac{d \ln f}{dv} = \frac{1}{w} \frac{d \ln f}{dw} = -\gamma$$

because it is probability...
$$f(u) = \sqrt{\frac{\gamma}{2\pi}} e^{-\frac{1}{2} n u^2}$$

Average kinetic energy (expectated value):

$$\left\langle \frac{1}{2} m u^2 \right\rangle = \int_{-\infty}^{+\infty} \frac{1}{2} m u^2 f(u) du = \int_{-\infty}^{+\infty} \frac{1}{2} m u^2 \sqrt{\frac{\gamma}{2\pi}} e^{-\frac{1}{2} \mu u^2} du = \frac{1}{2} m \sqrt{\frac{\gamma}{2\pi}} \int_{-\infty}^{+\infty} u^2 e^{-\frac{1}{2} \mu u^2} du$$

$$= \frac{1}{2}m\sqrt{\frac{\gamma}{2\pi}}\sqrt{\frac{2\pi}{\gamma^3}} = \frac{m}{2\gamma}$$
 using
$$= \frac{1}{2}kT$$
 therefor
$$\gamma = \frac{m}{kT}$$
 eq. of state

Maxwell distribution:

$$f(u) = \sqrt{\frac{m}{2\pi kT}}e^{-\frac{m}{2}u^2/kT}$$

$$f(u) = \sqrt{\frac{m}{2\pi kT}} e^{-\frac{m}{2}u^2/kT}$$

$$F(u, v, w) = \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2}(u^2 + v^2 + w^2)/kT}$$

Example: Ideal gas (3)

Maxwell distribution

$$F(u,v,w) = \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2}\left(u^2 + v^2 + w^2\right)/kT}$$

Probability function $\propto e^{\{\text{kineticenergy}\}/kT}$

$$\propto e^{\{ ext{kineticenergy}\}/kT}$$

$$f(u) = \sqrt{\frac{m}{2\pi kT}}e^{-\frac{m}{2}u^2/kT}$$



- Only for Ideal gas?
- < Remember >
- Empirical knowledge is used.

Eq. of state for ideal gas



Relation between Temp. and average Kinetic energy



This is the definition of Temp. !?

Let's consider more general case!

You may have learnt Thermodynamics

Thermodynamics

 $(0^{th} low : A.eq.B and B.eq.C -> A.eq.C)$

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Everything is derived from these principles.

- Definitions
 - Temperature
 - Free energy
 - Entropy
 - equilibrium
 - etc...

Temperature is kinetic energy?

(at equilibrium)

Temperature and equilibrium (general treatment)

Canonical ensemble (1)

Num. of elements with energy e_i is n_i . Total num. of elements (particle, molecule, etc...) : $N = \sum n_i$ N and E are given.

Why!?

..., therefore

 $n_i(e_i)$ is determined by the state. (density of state) States with same $n_i(e_i)$ are indistinguishable.

Assumption 2 (most important !!)

The state is observed so as to maximize the way of shuffling of the density of state $n_i(e_i)$

Canonical ensemble (2)

the way of shuffling of the state

 $W = \frac{N!}{n_1! n_2! n_3! \Lambda}$

To maximize W

$$L \equiv \ln W - \lambda_1 \left(\sum_i n_i - N \right) - \lambda_2 \left(\sum_i n_i e_i - E \right) \rightarrow \max$$

Lagrange multiplier

using the definition of W

$$L = \left(\ln N! - \sum_{i} n_{i}!\right) - \lambda_{1} \left(\sum_{i} n_{i} - N\right) - \lambda_{2} \left(\sum_{i} n_{i} e_{i} - E\right)$$

$$\frac{\partial L}{\partial n_i} = \frac{\partial L}{\partial \lambda_i} = 0$$

using the Stirling's formula $\ln x! \approx x \ln x - x$

$$L = \left(N \ln N - N - \sum_{i} \left(n_{i} \ln n_{i} - n_{i}\right)\right) - \lambda_{1} \left(\sum_{i} n_{i} - N\right) - \lambda_{2} \left(\sum_{i} n_{i} e_{i} - E\right) \rightarrow \max$$

Canonical ensemble (3)

$$L = \left(N \ln N - N - \sum_{i} \left(n_{i} \ln n_{i} - n_{i}\right)\right) - \lambda_{1} \left(\sum_{i} n_{i} - N\right) - \lambda_{2} \left(\sum_{i} n_{i} e_{i} - E\right) \rightarrow \max$$

Here, N and E are constant, therefore

$$\frac{\partial L}{\partial n_i} = -\left(\ln n_i + 1 - 1\right) - \lambda_1 - \lambda_2 e_i = 0$$

$$n_i = \exp\left(-\lambda_1 - \lambda_2 e_i\right) = \Lambda_1 \exp\left(-\lambda_2 e_i\right)$$



$$n_i = \exp(-\lambda_1 - \lambda_2 e_i) = \Lambda_1 \exp(-\lambda_2 e_i)$$

Using $N = \sum n_i$

$$\frac{n_i}{N} = \frac{1}{\sum_{i} \exp(-\lambda_2 e_i)} \exp(-\lambda_2 e_i) \propto \exp(-\lambda_2 e_i)$$

Comparing with Maxwell distribution $f(u) = \sqrt{\frac{m}{2\pi kT}}e^{-\frac{m}{2}u^2/kT}$

$$f(u) = \sqrt{\frac{m}{2\pi kT}} e^{-\frac{m}{2}u^2/kT}$$

$$\frac{n_i}{N} = \frac{1}{\sum_{i} \exp(-e_i/kT)} \exp(-e_i/kT)$$
: Maxwell-Boltzmann distribution

 e_i is not limited to be the kinetic energy. kinetic energy, potential energy, etc...

Canonical ensemble (3)

$$L = \left(N \ln N - N - \sum_{i} \left(n_{i} \ln n_{i} - n_{i}\right)\right) - \lambda_{1} \left(\sum_{i} n_{i} - N\right) - \lambda_{2} \left(\sum_{i} n_{i} e_{i} - E\right) \rightarrow \max$$

Here, N and E are constant, therefore

$$\frac{\partial L}{\partial n} = -(1) \exp(-\lambda_2 e_i)$$

Using

$$\frac{n_i}{N} = \frac{1}{\sum_{i=1}^{N} \epsilon_i}$$

Temperature is energy! (used in physics) in the most probable situation.

Comparing with

$$\frac{n_i}{N} = \frac{1}{\sum_{i} \exp(-e_i/kT)}$$
 : Maxwell-Boltzmann distribution

I distribution $f(u) = \sqrt{\frac{m}{2\pi kT}}e^{-\frac{m}{2}u^2/kT}$

 e_i is not limited to be the kinetic energy. kinetic energy, potential energy, etc...

Canonical ensemble (4)

- Probability of the state of energy E:
- $\propto \exp(-E/kT)$
- Probability of the state whose energy > E: $\propto \exp(-E/kT)$
- "Assumption 2"

The state is observed so as to maximize the way of shuffling of the density of state

The most probable situation

Equilibrium!

Equilibrium is,,, the most probable state !!

You may have learnt Thermodynamics

Thermodynamics

(0^{th} low : A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Everything is derived from these principles.

- Definitions
 - Temperature
 - Free energy
 - Entropy
 - equilibrium
 - etc...

Temperature is energy! (at equilibrium)

Equilibrium is the most probable state.

Definitions

Temperature

T. can be shifted and scaled, so as proportional to energy!

Free energy

Helmholtz F = U - TS

Gibbs G = U + PV - TS

Free energy is same as the energy in physics?

Why they need to be defined as above?

Entropy

What is this?

Equilibrium

The most probable state! So, we can observe always...

Free energy and entropy

Prediction of phenomena

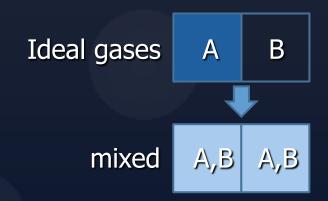
Physicists (kinetics)

as to decrease Potential Energy

m mgh

Chemists (chemistry)

as to decrease Free Energy



Decreased
Not changed
? (Not changed)

After the phenomena Potential Energy Total energy Free Energy

Not changed Not changed Decreased

Chemists wanted to introduce something like potential energy, to predict the phenomena.

Chemists wanted to introduce something like potential energy, to predict the phenomena.

The something; related to the work or energy, used in thermodynamic technology, (ex. steam engine, chemical reaction, etc...)

Definition of free energy:

Energy that can be converted into a work in a process of

constant temperature and volume (Helmholtz). F = U - TS

constant temperature and pressure (Gibbs).

G = U + PV - TS

(Two of three T, P, & V determine its state.)

Energy or work in kinetics

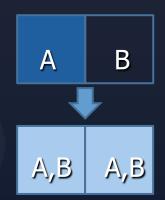
What's this!??

$$G = U + PV - TS$$

Definition of free energy:

Energy that can be converted into a work in a process





- The TS term increases in the mixing process, even if kinetic and potential energy never changes.
- The mixing process never goes back.

Increase in TS term \rightarrow Decrease of G

- = Increase of the usable energy
- = Degrade of energy

Definition of free energy:

Energy that can be converted into a work in a process

$$G = U + PV - TS$$

What's this!??

- There exists a part of energy, which we can not take as a work.
- It is proportional to the temperature.

The degrade is proportional to the temp...

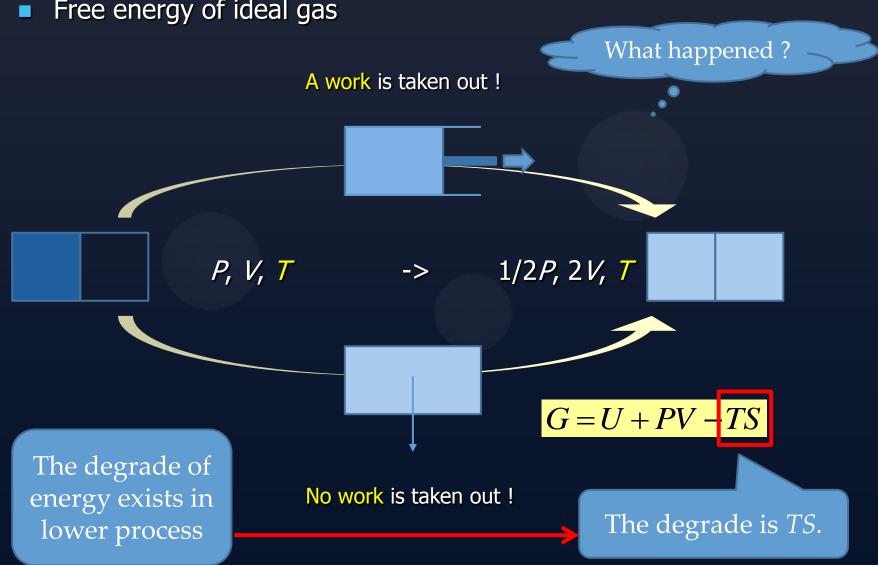
Degrade of the energy exits...

S is defined as the entropy. Empirical knowledge of chemists

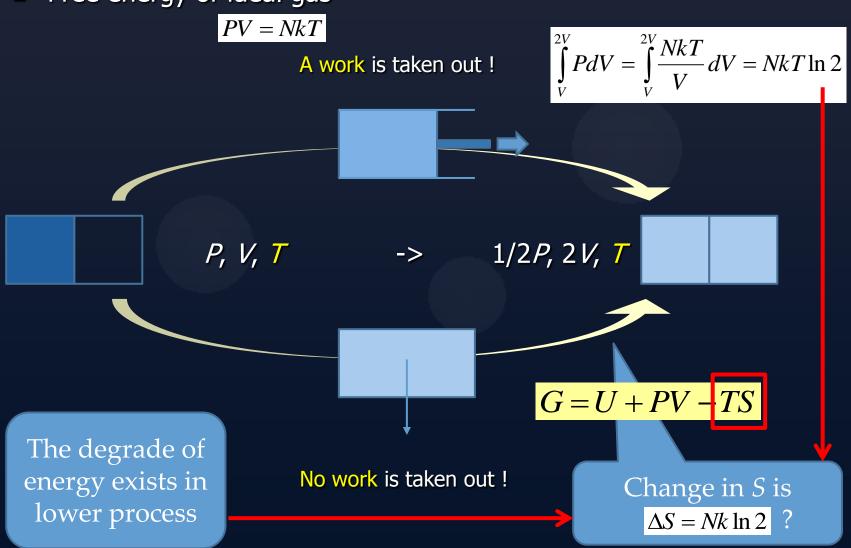
What's the entropy!??

This is the definition of entropy in thermodynamics.

Free energy of ideal gas



Free energy of ideal gas



Statistical Mechanics and Thermodynamics

Thermodynamics

(0^{th} low: A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Boltzmann's entropy formula:

 $S = k \ln W$



Grave-stone of Boltzmann...

Statistical Mechanics

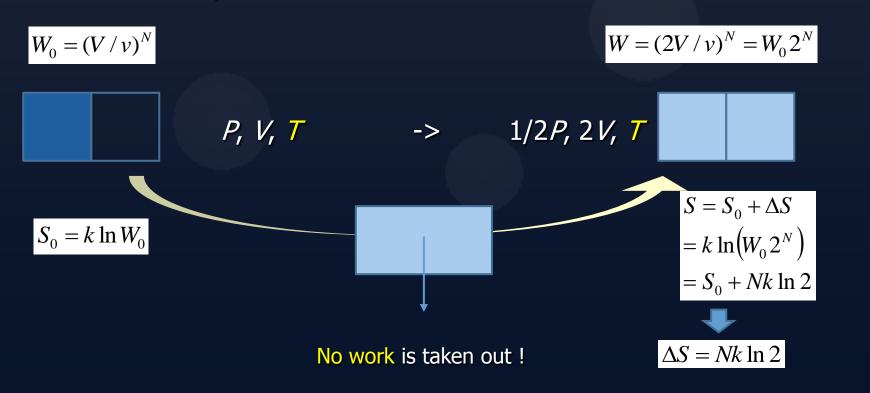
Ergodic theorem: time average = space average, (in equilibrium)

• W: Number of state $S = k \ln W$

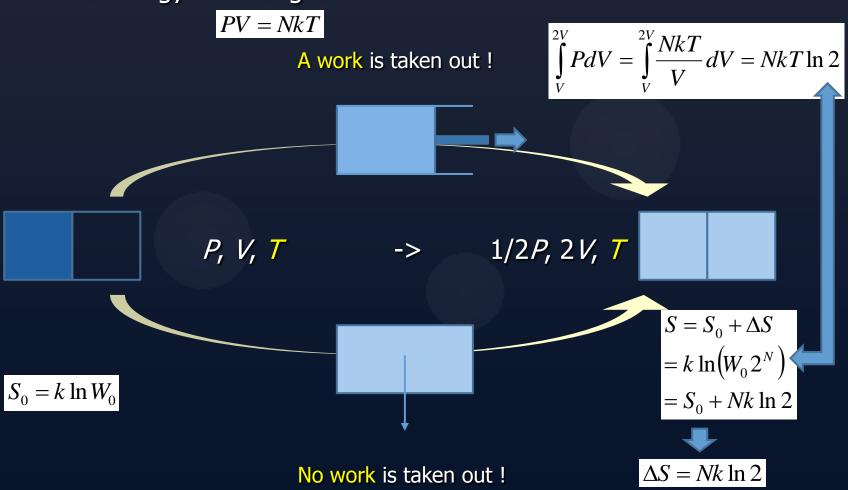
v: small control volume

V/v: num. of the c. volume 2V/v: num. of the c. volume

Num. of the way to shuffle to distribute the atoms to the each c. volume



Free energy of ideal gas



Free energy of ideal gas

Internal E. U_0

No change.

Free E.
$$G_0 = U_0 + PV - TS_0 ->$$

 G_0 - $T\Delta S$

Only entropy term must be changed as

 S_0

Entropy

 $S_0 + \Delta S$

must be

$$\Delta S = Nk \ln 2$$

Entropy change

Energy taken out : $\int_{0}^{2V} PdV = \int_{0}^{2V} \frac{NkT}{V} dV = NkT \ln 2$

Definition of entropy: $S = k \ln W$

If the entropy is defined as this, it becomes consistent.

$$S_0 = k \ln W_0$$

 $S_0 + \Delta S = k \ln(W_0 2^N) = S_0 + Nk \ln 2$

Check of understanding

Free energy of ideal gas

$$PV = NkT$$

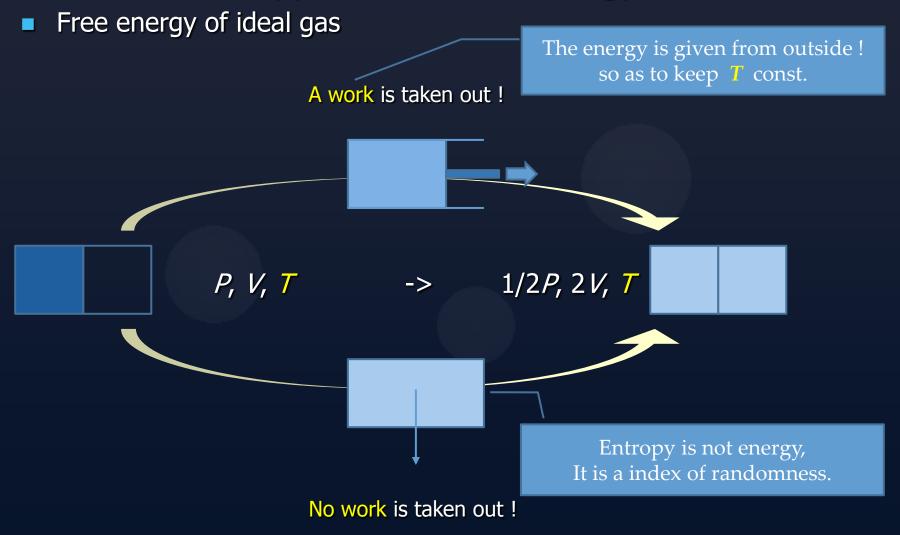
A work is taken out!

$$\int_{V}^{2V} PdV = \int_{V}^{2V} \frac{NkT}{V} dV = NkT \ln 2$$

We can take out the infinite work!???... although internal energy must be finite.

W depends on volume!???...
Energy never depend on volume...

$$S_0 = k \ln W_0$$
 No work is taken out
$$S_0 + \Delta S = k \ln (W_0 2^N) = S_0 + Nk \ln 2$$



Introduction of Free energy & entropy & temperature

Definition of Free Energy :

Energy that can be converted into a work.

(There exist the energy part we can not take out.)

Internal energy = Kinetic energy

we can take out from outside (by heat transfer...)

$$G = U + PV - TS$$

Degrade of the energy...

Definition of EntropyRelated to the randomness

$$S = k \ln W$$

← Definition of Temperature

Temperature

Empirical temp. °C : water defined one or two temp. °F : salt water

T. can be shifted and scaled, so as proportional to energy!

Free energy

Helmholtz F = U - TSGibbs G = U + PV - TS

Free energy is energy that can be converted into work. If defined as above, it become consistent.

Entropy

 $S = k \ln W$ corresponding to number of states, i.e. randomness.

Equilibrium

Most probable state! So, we can observe always...

These are the interface between Thermodynamics and statistical mechanics

Statistical Mechanics and Thermodynamics

Thermodynamics

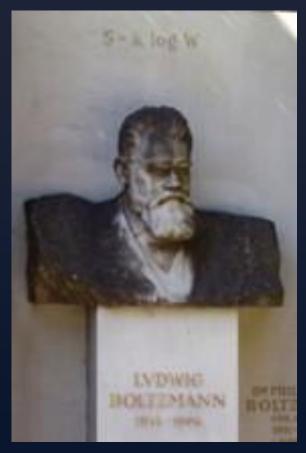
(0^{th} low: A.eq.B and B.eq.C -> A.eq.C)

1st low: Energy conservation + ...

2nd low: Free energy decreases

3rd low: Entropy -> 0 at 0K

Boltzmann's entropy formula : $S = k \ln W$



Grave-stone of Boltzmann...

Statistical Mechanics

Ergodic theorem: time average = space average, (in equilibrium)

Based of statistical mechanics, Thermodynamics can be utilized!

Activation energy and Arrhenius plots and rate-determining step

chemical reaction

diffusion coefficient

others... (evaporation, adsorption, etc...)

Arrhenius plots is to know the mechanism of rate-determining step.

Specific heat capacity (1: ideal gasses)

Remember thermodynamics! (How have you learnt?)

$$C_V = \frac{1}{N} \left(\frac{\partial U}{\partial T} \right)_V$$
 why? $C_p = \frac{1}{N} \left(\frac{\partial H}{\partial T} \right)_p$ why?

- Definition
 - Heat required to increase unit temp.
 - under constant volume

$$\Delta Q = \Delta G = \Delta U = NC_V \Delta T$$

- under constant pressure

$$\Delta Q = \Delta G = \Delta U + P\Delta V - \Delta T\Delta S = \Delta U + P\Delta V = NC_P\Delta T$$



$$C_p - C_V = k$$







$$G = U + PV - TS$$

Specific heat capacity (2: metals)

- Einstein's model
- Dulong-Petit's Low

What is the *U*?

- Debye model
 - Ionic cores

$$U = \langle e_{k} \rangle + \langle e_{p} \rangle = 3\frac{1}{2}kT + 3\frac{1}{2}kT = 3kT$$

- Electrons = negligible

$$\frac{\Delta Q}{\Delta T} = 3k$$

Quantum statistical mechanics

Phase diagram

- Stability of phases
- Phase rule
- Phase diagram
- Calculation of Phase Diagram (CALPHAD)